

PRECISE MEASUREMENT METHOD FOR TEMPERATURE COEFFICIENT  
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## ABSTRACT

A precise measurement method for the temperature coefficient of dielectric resonator material was developed. The error of measurement was decreased from 0.5 ppm/°C to 0.05 ppm/°C compared with the conventional method. This measurement method made techniques for improving stability possible, such as for filters, multiplexers and oscillators.

## INTRODUCTION

For filters and oscillators, a precise measurement technique for temperature coefficient has been needed in order to achieve improved temperature stability [1]. Recently, dielectric resonator materials with extremely small dielectric loss and better, more stable temperature coefficient were developed [2]. The requirement for a precise method of measuring the temperature coefficient was satisfied to the level of  $\pm 0.05$  ppm/°C. Also, because the linearity of the temperature coefficient of the dielectric constant  $K$  is one of important factors for achieving high temperature stability, there has been an added impetus to creating a precise method for temperature coefficient measurement.

Conventionally, temperature coefficients of dielectric resonator materials were measured with a dielectric rod resonator TE<sub>011</sub> mode, short-circuited at both ends by two parallel conducting plates [3]. However, this method is not fit for measurement of improved dielectric loss material. The precision of this type of measurement decreases because the unloaded Q of this measuring method is low compared with the Q factor of the material, and because the dimensional accuracy required is so high that it is impossible to achieve.

In this paper, we propose a precise measurement method using a TE<sub>011</sub> mode dielectric resonator. The precision of this temperature coefficient measuring method can satisfy the requirement of  $\pm 0.05$  ppm/°C because the unloaded Q of the measurement is as high as the Q factor of the dielectric materials.

## CONSTRUCTION OF MEASURING INSTRUMENT

A cross section of the measuring instrument is shown in Figure 1. The dielectric resonator is fixed to the support in the center of the cavity.

Figure 1-a shows the measuring instrument of a brass metal cavity of which the thermal expansion coefficient,  $\alpha$  is 18 ppm/°C. An Invar metal cavity of which  $\alpha$  is 0 ppm/°C is also used.

The shielding cavity, shown in Figure 1-b, are made of metalized ceramic material having the same  $\alpha$  as the dielectric resonator. The support is made of a ceramic

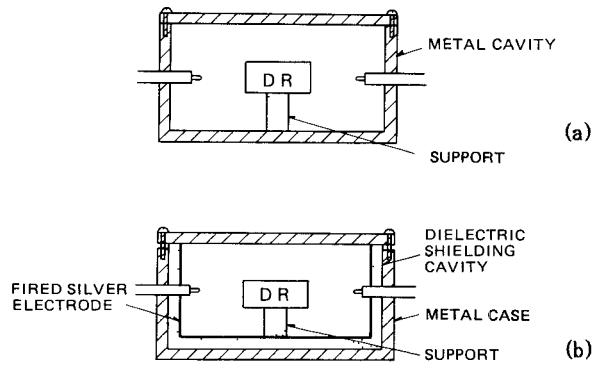


Figure 1 The cross sectional view of measuring instrument : (a) metal cavity  
(b) ceramic cavity

tube which also has the same  $\alpha$ .  $\alpha$  of the cavity and support material (2MgO·SiO<sub>2</sub>-ZrSiO<sub>4</sub>) can be controlled from between 4 ppm/°C to 10 ppm/°C using a composition ratio of 2MgO·SiO<sub>2</sub> and ZrSiO<sub>4</sub>.  $\alpha$  of the ceramic material (2MgO·SiO<sub>2</sub>-ZrSiO<sub>4</sub>) is shown in Figure 2.

The difference of the thermal expansion caused by  $\alpha$  of the dielectric resonator and the cavity is shown in Figure 3. The dimension ratio of the dielectric resonator and the metal cavity dimension was changed by the thermal expansion; however, this ratio, in the case of a ceramic cavity, was not changed by thermal expansion. Therefore, using a ceramic cavity, the temperature coefficient of the material can be measured without the effect of the thermal expansion.

## PRINCIPLES

We considered the measurement of the temperature dependence of resonant frequency. When the measuring instrument of Figure 1-b is used, the normalized resonant frequency is that shown in the following equation,

$$\frac{2\pi D}{\lambda_0} = F(K_1, K_2, L/D, d_1/D, d_2/D, \ell_1/D, \ell_2/D) \quad (1)$$

where,  $\lambda_0$  is resonant wavelength in free space,  $K_1$  and  $K_2$  are dielectric constant of a resonator and a support. If the small linear changes  $\Delta D$ ,  $\Delta K_1$ ,  $\Delta K_2$ ,  $\Delta L$ ,  $\Delta d_1$ ,  $\Delta d_2$ ,  $\Delta \ell_1$ ,  $\Delta \ell_2$ ,  $\Delta f$  in (1) are caused by the temperature change  $\Delta T$  of the instrument, the following equation is derived.

$$\eta_f = \frac{1}{F} \frac{\partial F}{\partial K_1 \Delta T} + \frac{1}{F} \frac{\partial F}{\partial K_2 \Delta T} + \frac{1}{F} \sum_i \frac{\partial F}{\partial X_i \Delta T} - \frac{1}{D} \frac{\Delta D}{\Delta T} \quad (2)$$

where

$$\eta_f = \frac{1}{f} \frac{\Delta f}{\Delta T}, X_i = L/D, d_1/D, d_2/D, \ell_1/D, \ell_2/D \quad (3)$$

Here,  $\eta_f$  is the temperature coefficient of resonant frequency. If  $\alpha$  of a resonator, a cavity, and a support are all equal, then  $\Delta X_i = 0$  in (2). Therefore the following equation is derived from (2).

$$\eta_f = \frac{1}{F} \frac{\partial F}{\partial K_1 \Delta T} + \frac{1}{F} \frac{\partial F}{\partial K_2 \Delta T} - \alpha \quad (4)$$

where

$$\alpha = \frac{1}{D} \frac{\Delta D}{\Delta T} \quad (5)$$

$\frac{1}{F} \frac{\partial F}{\partial K_i \Delta T}$  is the change ratio of resonant frequency which is caused by the temperature coefficient of  $K$ . The change ratio is shown in the following equation [4].

$$\frac{1}{f_0} \frac{\Delta f}{\Delta T} = -\frac{1}{2} A \frac{1}{K} \frac{\Delta K}{\Delta T} \quad (6)$$

where,  $A$  is the ratio of the electric energy stored in the dielectric rod in proportion to the total electric energy in the cavity.

The following equation is derived from (4), (6).

$$\eta_f = -\frac{1}{2} (A_1 \eta_{K_1} + A_2 \eta_{K_2}) - \alpha \quad (7)$$

where

$$\eta_{K_i} = \frac{1}{K_i} \frac{\Delta K_i}{\Delta T} \quad (8)$$

If all electric energy is stored in the dielectric rod,  $A_1=1$  and  $A_2=0$  in (7).  $\eta_{f0}$  defined when  $A_1=1$  and  $A_2=0$ , is shown in the following equation.

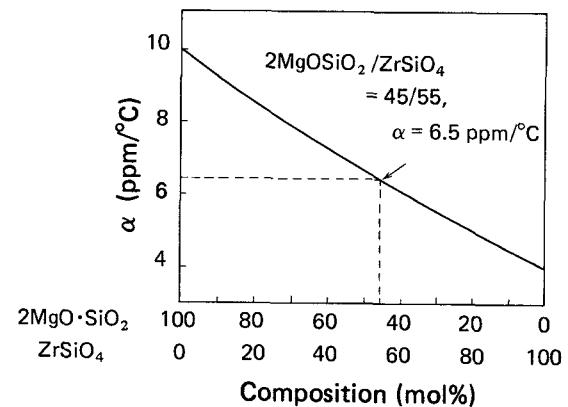


Figure 2 The thermal expansion coefficient of the ceramic material

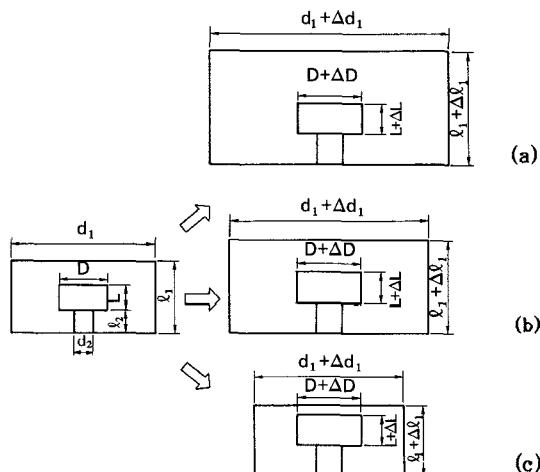


Figure 3 The thermal expansion of the dielectric resonator and the cavity : (a) brass cavity, (b) ceramic cavity, (c) Invar cavity

$$\eta_{f_0} = -\frac{1}{2} \eta_{K_1} - \alpha \quad (9)$$

Finally, the temperature coefficient of frequency is given by the following equation derived from (7), (9).

$$\eta_f = \eta_{f_0} - \frac{1}{2} (1 - A_1) \eta_{K_1} + \frac{1}{2} A_2 \eta_{K_2} \quad (10)$$

### MEASURED RESULTS

The measuring system is shown in Figure 4. Temperature coefficient for the dielectric resonator of  $K = 37.6$ ,  $D = 10.1 \text{ mm}$  and  $L = 4.9 \text{ mm}$  was measured by using the Invar cavity, the ceramic cavity, and the brass cavity.  $\Delta\eta_f$ , the measured value deviation of  $\eta_f$  from  $\eta_{f_0}$  is shown in Table 1. ; its mean-square errors are about 0.05 ppm/°C. The histogram of measured  $\Delta\eta_f$  is shown in Figure 5.  $\Delta\eta_f$  in Table 1 is measured with the dielectric rod in the center height of the cavity; however,  $\Delta\eta_f$  has dependence on the resonator position in the metal cavity because of the difference between the thermal expansion of the ceramics and the metal. The  $\Delta\eta_f$  dependence on the resonator position is shown in Figure 6. By increasing the support length,  $\Delta\eta_f$  measured in the Invar cavity increased and  $\Delta\eta_f$  measured in the brass cavity decreased.  $\Delta\eta_f$  can be measured without the dependence on the resonator position in the ceramic cavity.

The temperature coefficient of the three materials ( $K=30.0, 37.6, 89.0$ ) were measured using the conventional method and this newly developed method. The measured results are shown in Table 2. The measured unloaded  $Q$ ,  $Q_0$  of the conventional method are higher than the one of this newly developed method. The mean-square error is decreased from 0.5 ppm/°C to 0.05 ppm/°C compared with the conventional method.

Table 1. Measured  $\Delta\eta_f$  of three cavities

Cavity material	Invar	Ceramic	Brass
$\alpha$ of cavity (ppm/°C)	0.0	6.5	18.0
Measured $\Delta\eta_f$ (ppm/°C)	$-0.1 \pm 0.05$	$-0.1 \pm 0.05$	$-1.0 \pm 0.05$

#### (Measuring condition)

Resonator :  $D=10.1 \text{ mm}$ ,  $L=4.9 \text{ mm}$ ,  $K=37.6$ ,  $\eta_{f_0}=0.25 \text{ ppm/}^{\circ}\text{C}$

Cavity :  $d_1=32.3 \text{ mm}$ ,  $l_1=14.3 \text{ mm}$

Support :  $d_2=5.0 \text{ mm}$ ,  $l_2=4.8 \text{ mm}$ ,  $K=8.5$ ,

$\eta_f=-60 \text{ ppm/}^{\circ}\text{C}$

Frequency :  $f_0=5.1 \text{ GHz}$

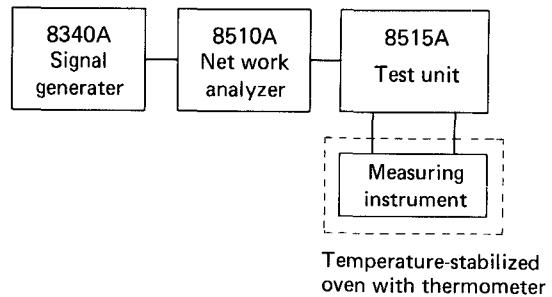


Figure 4 The measuring system

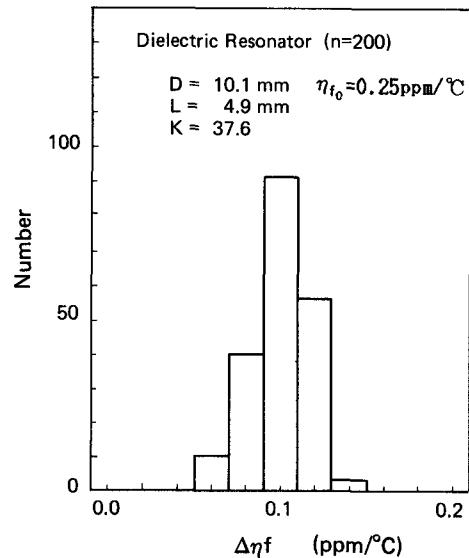


Figure 5 The histogram of measured  $\Delta\eta_f$

Table 2. Measured  $\eta_f$  of three resonators

Resonator material	A	B	C
$K$ of resonator	30.0	37.6	89.0
$\alpha$ of resonator (ppm/°C)	10.5	6.5	8.5
Conventional method:			
$\eta_f$ (ppm/°C)	$2.00 \pm 0.50$	$0.10 \pm 0.50$	$-5.00 \pm 0.50$
Frequency (Q <sub>0</sub> )	6.0 (3500)	5.3 (3000)	3.5 (1200)
Developed method:			
$\eta_f$ (ppm/°C)	$2.10 \pm 0.05$	$0.15 \pm 0.05$	$-4.92 \pm 0.05$
Frequency (Q <sub>0</sub> )	5.8 (12000)	5.1 (9000)	3.4 (2000)

#### (Measuring condition)

Resonator:  $D=10.1 \text{ mm}$ ,  $L=4.9 \text{ mm}$

## CONCLUSION

The precise measurement method for temperature coefficient of dielectric resonator materials was developed. The measuring instrument is shown in Figures 7 and 8.

This method is suitable for measuring the temperature coefficient of the low-loss dielectric material for high selective filters [5].

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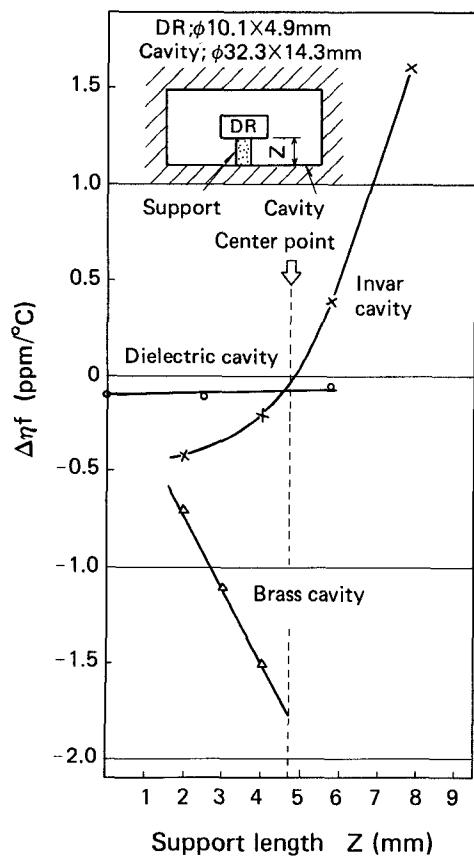


Figure 6 The measured results of  $\Delta\eta_f$

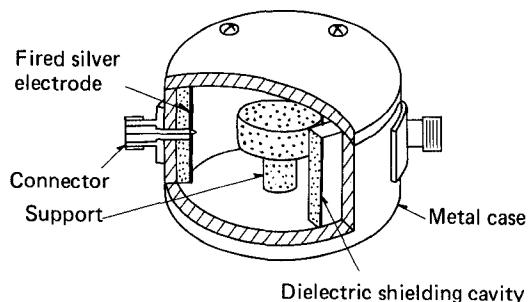


Figure 7 The measuring instrument

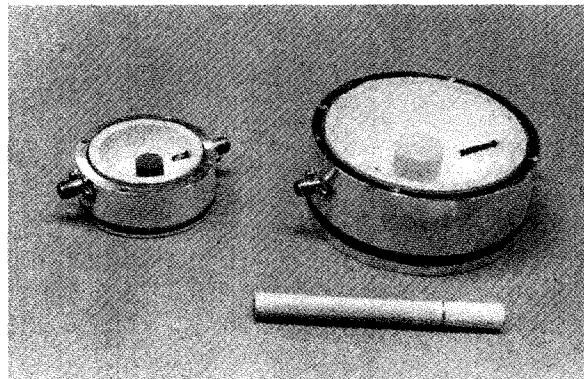


Figure 8 The photograph of the measuring instrument